



Satellites: Jacobian Integral of the equation of motion and elliptic orbit

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Abstract- This paper deals with the elliptic motion of center of mass, periodic Oscillation and it provides stable solution to the equation of oscillations. The change of frequencies of oscillation doesn't affect the stability of the Solution.

Keywords: stability; periodic oscillation; frequencies

1. INTRODUCTION

The effect of atmosphere resistance magnetic force and oblateness of the earth on the motion and stability of two satellites connected in the particular case of keplerian orbit that is circular orbit. The system moving in space has to face same perturbative force. These Perturbative forces compel the system which changes its orbit from circular to elliptic one.

2. MATHEMATICAL APPROCH

This system of equations of motion may be written in the following form

$$\begin{aligned} x'' - 2y' - 3x\rho + \frac{4A}{\rho}x + f\rho\rho' + \frac{B}{\rho}\cos i &= \\ -\lambda\bar{\alpha}\left[1 - \frac{\ell o}{\rho_r}\right]x\rho^4 & \\ y'' + 2x' + f\rho^2 - \frac{A}{\rho}y + \frac{B\rho'}{\rho^2}\cos i &= \\ -\lambda\bar{\alpha}\rho^4\left[1 - \frac{\ell o}{\rho_r}\right]y & \\ z'' + z - \frac{A}{\rho}z = & \\ -\lambda\bar{\alpha}\rho^4\left[1 - \frac{\ell o}{\rho_r}\right]z\frac{-B}{\rho}\left[\frac{\rho'}{\rho}\cos(v+\omega) + \frac{1}{\mu_E}(3\rho^3\rho^2 - \right. & \\ \left. \mu_E)\sin(v+\omega)\right] \sin i & \quad \text{----- (1)} \end{aligned}$$

Where,

$$r = \sqrt{x^2 + y^2 + z^2}$$

In case of elliptic orbit of the center of mass the dashes denote differentiation with respect to v of the orbit.

The condition of constraint takes the form

$$x^2 + y^2 + z^2 \leq \frac{\ell o^2}{\rho^2} \quad \text{----- (2)}$$

The equations (1) are non-homogeneous and non-linear in character with the periodic forces .The appearance of the periodic effects on the motion of the satellite m_1 so long the period of oscillation of the system without their periodic force.

We assume that there periods are non-commensurate the equations of motion can be averaged with respect to v and the averaged equations of motion will describe the small perturbations and long periodic

effects due to the periodic force on the motion of the satellite m_1 . Now, we may obtain the averaged values the following terms :

$$\rho = \frac{1}{1+e\cos v}$$

$$\begin{aligned} \frac{1}{2\pi} \int_0^{2\pi} \rho dv &= \frac{1}{(1-e^2)^{1/2}} \\ \frac{1}{2\pi} \int_0^{2\pi} \rho \rho' dv &= 0 \quad \frac{1}{2A} \int_0^{2\pi} \frac{\rho^1}{\rho^2} dv \\ \frac{1}{2\pi} \int_0^{2\pi} \rho^2 dv &= \frac{1}{(1-e^2)^{3/2}} \end{aligned}$$

$$\frac{1}{2\pi} \int_0^{2\pi} \rho^3 dv = \frac{(2+e^2)}{2(1-e^2)^{5/2}} \quad 3$$

$$\begin{aligned} \frac{1}{2\pi} \int_0^{2\pi} \rho^4 dv &= \frac{(2+3e^2)}{2(1-e^2)^{7/2}} \\ \frac{1}{2\pi} \int_0^{2\pi} \frac{\rho^1}{\rho^2} \cos(v+\omega) dv &= 0 \\ \frac{1}{2\pi} \int_0^{2\pi} \rho^2 \sin d(v+\omega) v &= 0 \\ \frac{1}{2\pi} \int_0^{2\pi} \sin d(v+\omega) v &= 0 \end{aligned}$$

Here,

$$\begin{aligned} x'' - 2y' - \frac{3}{(1-e^2)^{1/2}}x + 4Ax + B\cos i &= \\ = -\lambda\bar{\alpha}\left[\frac{(2+3e^2)}{2(1-e^2)^{7/2}} \right. & \\ \left. - \frac{(2+e^2)\ell o}{2r(1-e^2)^{5/2}}\right]x & \end{aligned}$$

$$\begin{aligned} y'' + 2x' + \frac{f}{(1-e^2)^{3/2}} - A.y &= \\ = -\lambda\bar{\alpha}\left[\frac{(2+3e^2)}{2(1-e^2)^{7/2}} \right. & \\ \left. - \frac{(2+e^2)\ell o}{2r(1-e^2)^{5/2}}\right]y & \end{aligned}$$



$$z'' + z - A.z = -\lambda\bar{\alpha} \left[\frac{(2+3e^2)}{2(1-e^2)^{7/2}} - \frac{(2+e^2)\ell o}{2r(1-e^2)^{5/2}} \right] z$$

Where

$$r^2 = x^2 + y^2 + z^2 \quad \text{--- (4)}$$

∴ The system of equations given by (4) will represent the motion of the system while the relation

$$x^2 + y^2 + z^2 \leq \ell o^2(1-e^2)^{3/2} \quad \text{----- (5)}$$

Gives the condition of the constraint we can say that the satellite m_1 is moving on the spare given by

$$x^2 + y^2 + z^2 = \ell o^2(1-e^2)^{3/2} \quad \text{----- (6)}$$

The equation (4) must exist Jacobi's integral for the system of equation (4) multiplying the three equations by $2x'$, $2y'$, $2z'$ and adding them together them after integrating

$$\begin{aligned} x'^2 + y'^2 + z'^2 - \left[\frac{3}{(1-e^2)^{3/2}} - 4A \right] x^2 - A.y^2 \\ + z^2 - A.z^2 + 2Bxcosi \\ + \frac{2f}{(1-e^2)^{3/2}} y \\ + \lambda\bar{\alpha} \frac{(2+3e^2)}{2(1-e^2)^{7/2}} (x^2 + y^2 + z^2) \\ - \frac{\lambda\bar{\alpha}\ell o(2+e^2)}{(1-e^2)^{5/2}} [x^2 + y^2 + z^2]^{1/2} \\ + he \quad \text{--- (7)} \end{aligned}$$

We have obtained the Jacobi's integral in the form of the equation (7) and he will be treated as the Jacobian constant. By obtaining the equation of the surface of zero velocity of the system .It follows that the satellite m_1 will be moving within the boundary of the surfaces of zero velocity.

REFERENCES

- [1]. Beletsky ,V.V Motion of artificial satellite relative to the center of mass, Nauka, 1965(Russia).
- [2]. Jlatustev,V.A Okhotsinsky,D.E Sarichev V.A & Tefhevsky A.P, Periodic Solution of the problem of planar oscillation of satellite in elliptical orbit.
- [3]. Sarichev,V.A; Sazonov,V.V & Zlatustov, Periodic oscillation of a satellite in the plane of an elliptic orbit cosmic Research vol.15 P 698
- [4]. Singh R.B,The motion of two connected bodies in an elliptical orbit, Bulletin of the Moscow state University Mathematics Mechanics ,No,3PP,82-88,1973(Russia)
- [5]. Yu, E.Y.,Optimum design of a gravitationally oriented two body satellites. Bell system Technical Journal, Vol44 January 1965
- [6]. Austin, F. : Free non-linear dynamics of a rotating flexibly connected double mass station Rept. ADR 06-15-64, 1, Jan. 1964, Gunman Aircraft Engineering Corporation.
- [7]. Brouwer, D & Clemence, G.M. : Method of celestial mechanics Academic Press, New York, 1961
- [8]. Burnside, W.S. & Panton, A.W.; The theory of equation vol. 1, S. Chand and Company, New Delhi, 1979
- [9]. Chovotov, V. : Gravity gradient excitation of a rotating cable-counter –weight space station in orbit, Journal of applied Mechanics, Vol. 30, No. 4PP. 547-554 Dec. 1963 .
- [10]. Chetayev, N.G. : Stability of Motion 3rd Ed., "Nauka" Moscow, 1966 (Russian)
- [11]. Danby, J.M. : Fundamentals of celestial mechanics The Mcmillan Co., New York 1962
- [12]. Demin, V.G. : Motion of artificial satellite in non-central field of attraction, Nauka, Moscow, 1968 (Russian)
- [13]. Sharma, Amit, and Ashish Kumar Sharma. "Mathematical Modeling of Vibration on Parallelogram Plate with Non Homogeneity Effect." *Romanian Journal of Acoustics and Vibration* 13, no. 1 (2016): 53.
- [14]. Sharma, Subodh Kumar, and Ashish Kumar Sharma. "1651. Rayleigh-Ritz method for analyzing free vibration of orthotropic rectangular plate with 2D thickness and temperature variation." *Journal of Vibroengineering* 17, no. 4 (2015).
- [15]. Elsgotts, L. : Differential equations and the calculus of variations, PP. 211 – 250 Mir Publishurs Moscow, 1973
- [16]. Etkin, B.: Dynamics of gravity oriented orbiting systems with applications to passive stabilization. A.I.A.A. Journal, Vol.2, No. 6 PP. 1008 – 1014 June, 1964
- [17]. King, D. G. – Hele; Theory of satellite orbits in an atmosphere, Butterworth, London, 1964